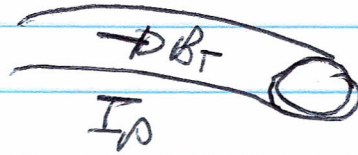


Physics 817Class 5

~ recall tokamak



~ toroidal pinch

~ strong field, $q \geq 1$, $q' > 0$

~ MHD stable

~ confinement (heat Bohm) $\frac{eT}{k_B} = W_B$

60's

and

~ auxiliary heating raised temperature

~ OH scaling established

$$T_E \sim n R^2 q$$

70's

~ partial pinch raised central n , peaked $n(r)$.

Physics

~ 80's, 90's = Decade of Confinement Studies, Improvement

- in Ohmic tokamak; T , current tightly linked

$$nT = E_T = \frac{V_L}{2\pi R} \rightarrow \text{const.}$$

$$n \sim 1/T^{3/2}$$

→ Aside: 80's → decade of confidence

80's saw the construction of large tokamaks

TFTR - Princeton (PPPL) $T_i \sim 40 \text{ keV}$
 → 1998

JET - Abingdon, still operating

JT-60 - Naka, Japan
 → 2008

JET: $R \sim 3 \text{ m}$, $a \sim 1.25 \text{ m}$

$B \sim 4 \text{ T}$

$I_p \sim 1 \text{ MA}$

$P \sim 25 \text{ MW}$

$T_i \sim 40 \text{ keV}$

$P_R \sim 1/6 \text{ MW}$

$Q \sim 6$

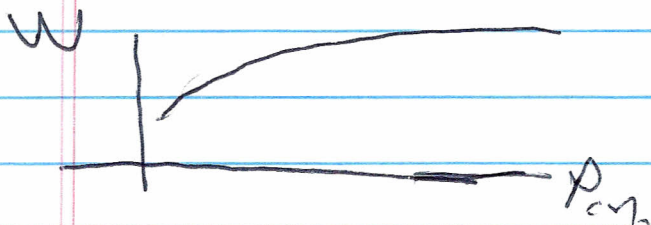
- with P_{aux} ,

- τ_E degrades with power

$$\tau_E \sim P^{-\alpha}$$

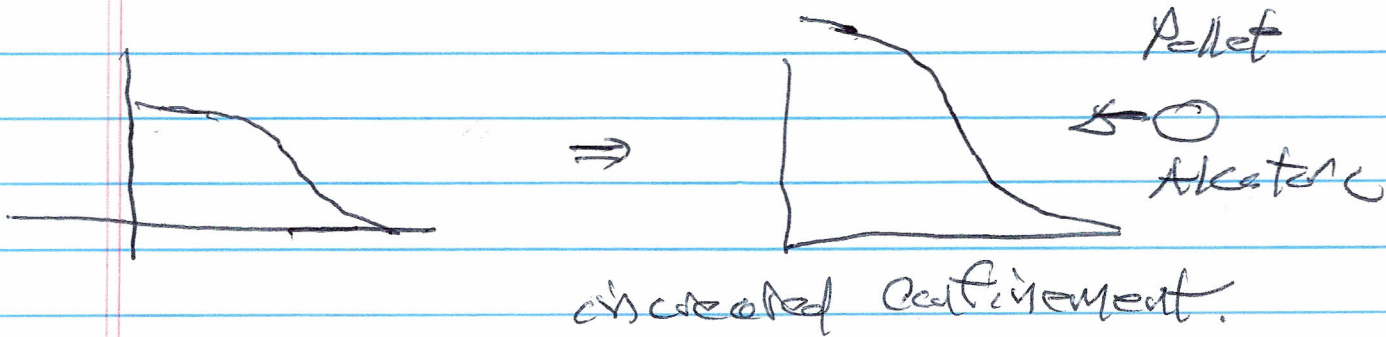
$$- \tau_E = \frac{W}{P_{in}}$$

\Rightarrow energy content saturates/exhibits diminishing returns to scale with P_{in}



Why? \Rightarrow transport mechanism?

- good news: peaked density profiles favorable

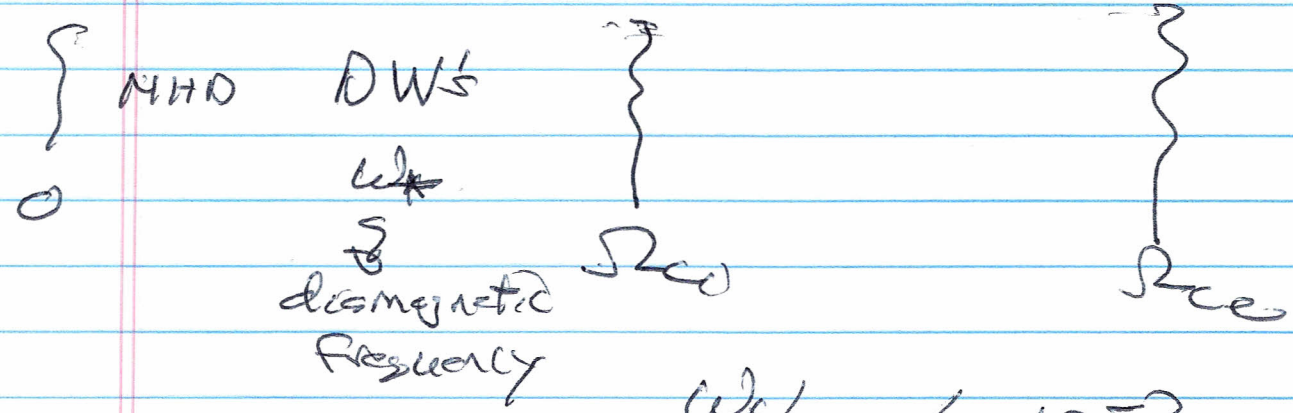


Why? - Drift Waves

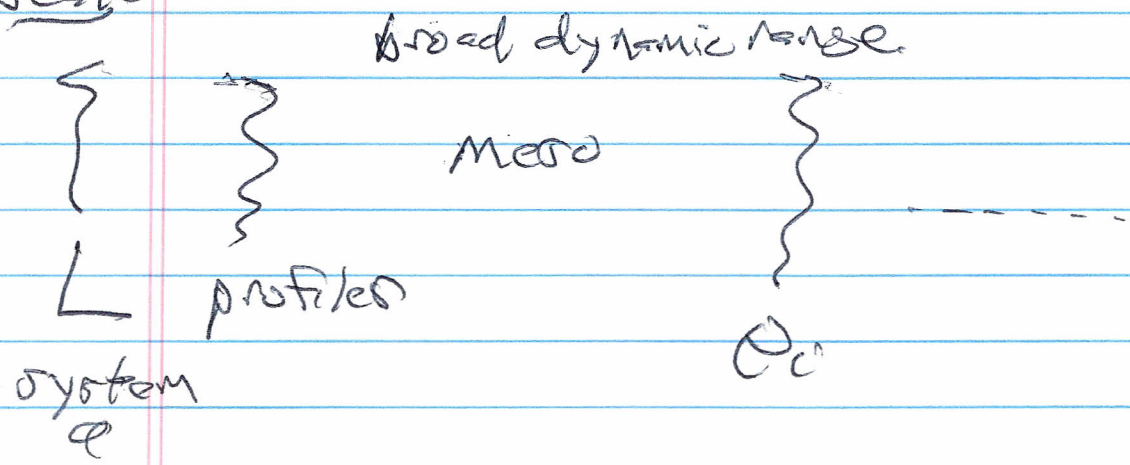
- Drift waves are the "microinstabilities" that drive energy particle transport

→ Scales

Frequency



Scale



Point $\lambda_p / L_p \lesssim 10^{-2}$

Study

→ predicted in former USSR in late 50's

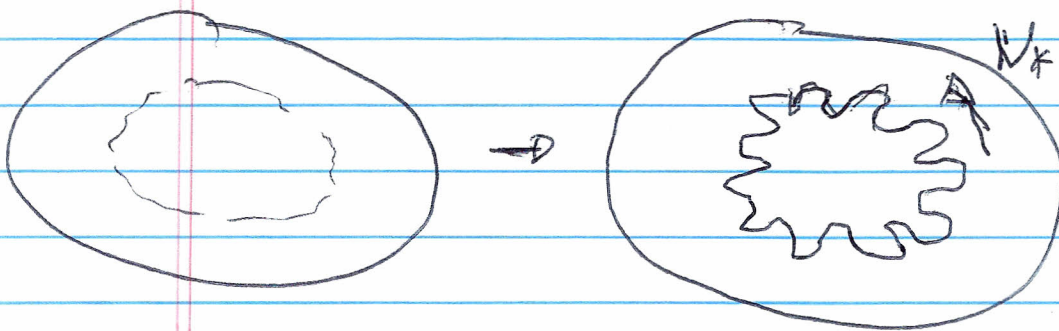
⇒ driven by focus on Bohm scaling of linear devices

→ observed in late 70's by $\left. \begin{matrix} \text{microwaves} \\ \text{laser} \end{matrix} \right\}$

scattering ⇒ detected density fluctuations

$(\delta n/n)^2_{k,\omega}$ in $\left[\begin{matrix} \text{relevant scale, } \omega \text{ range} \\ \text{relevant size} \end{matrix} \right]$

→ J. M. S. → some of the other



⊙ Plate



How fast : $V_f \rightarrow$ diamagnetic velocity

i.e. $\frac{\partial n}{\partial t} + \frac{V_{EXB}}{B} \cdot \nabla n_0$

$$\frac{n_e}{n_0} = \frac{e \mu_0 \Phi}{T} = \frac{n_e}{n_0}$$

$$-c\omega \frac{\delta n}{n_0} = + \frac{c}{B} \nabla \Phi \frac{\partial n_0}{\partial r}$$

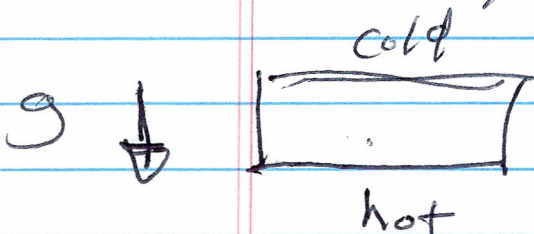
$$\omega = k_0 V_f$$

$$V_f = \frac{c}{L_n}$$

Unstable :

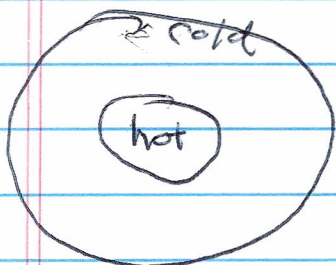
- e-i friction

- aka Rayleigh-Benard

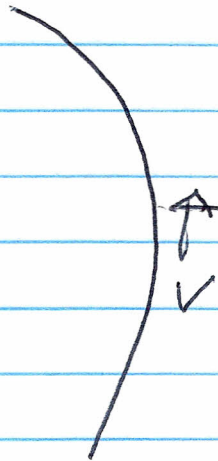


$$\Delta T > \Delta T_{crit}$$

\Rightarrow heat used by buoyancy



~~analogue~~ analogue of \underline{j} \Rightarrow magnetic curvature induced centrifugal force



$$\frac{F}{R} = mV^2$$

$$\frac{1}{R_c}$$

\hookrightarrow radius of curvature of field lines

\Rightarrow R-B micro-convection cells in tokamak.

Now, F-O-M : $D \rightarrow$ diffusivity

$$D \sim \lambda V \sim \lambda V_t$$

$$\sim \lambda \frac{\rho_s c_s}{L_n}$$

$$\sim \frac{\lambda}{L} D_B$$

what is $\lambda_I^2 \rightarrow$ few ρ_c

$$\Rightarrow D \sim \underbrace{\frac{\rho_e}{L} D_B}_{\text{Gyro-reduced Bohm}}$$

Gyro-reduced Bohm.

$$\sim \frac{T^{3/2}}{B^2}$$

- ρ_e/L smaller than Bohm
but

- $\sim T^{3/2} \rightarrow$ strong degradation

Whole can eliminate some but not all,
via peaked profiles, not a ~~good~~
prognosis...

\Rightarrow Once again, a surprise self-
organization intervened!

H-Mode

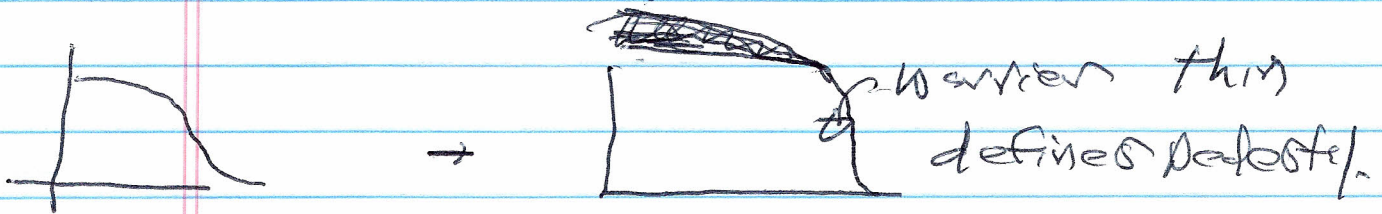
ASDEX Tokamak
Germany 1/1982
and many other
contributors.

L \rightarrow low

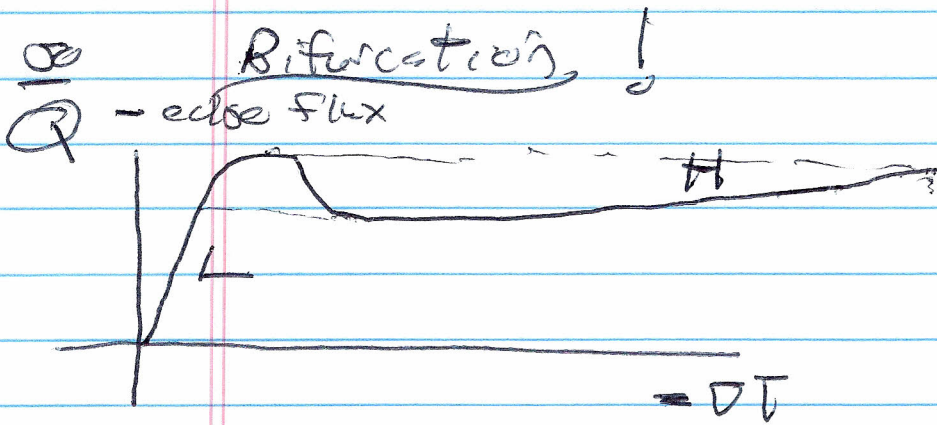
H \rightarrow high

→ W jumps up, for $P > P_{crit}$

→ A transport barrier / thermal insulation layer forms at edge



→ in barrier, fluctuations \ominus extinguished.



$$Q = -k \Delta T$$

→ mechanism seems to be onset of edge shear flow, ripping up eddies

